## The Effect of Spectral Height on Low-Frequency Motion Amplitudes

Simplified Treatment for Lightly Damped Linear Systems and Rectangular Spectra

We examine the qualitative effect on low-frequency motions of a moored vessel to changes in wave spectral amplitude by analyzing vessel response to two rectangular wave spectra,  $S_1$  and  $S_2$ , of equal variance. Waves associated with these spectra approach head-on to a vessel in a lightly damped linear mooring restraint. We assume for simplicity that regular-wave drift-force coefficients are frequency independent with constant value  $f(\omega) = f_0$ . The regular wave drift force  $F_{RW}$  is

$$F_{RW} = 0.5 * f_0 * D_w * B * \eta^2$$
.

Here,  $D_w$  is the weight density of water, B is the beam of the vessel and  $\eta$  is wave amplitude. Under the stated conditions these spectra posses the same variance  $(\sigma_w^2)$  and associated mean wave drift force  $(F_{MD})$ ; "U"pper and "L"ower rectangular frequency limits are  $\omega_U$ ,  $\omega_L$  respectively:

$$\sigma_{w}^{2} = S_{1} * (\omega_{U_{1}} - \omega_{L_{1}}) \equiv S_{1} * \Delta \omega_{1} = S_{2} * (\omega_{U_{2}} - \omega_{L_{2}}) \equiv S_{2} * \Delta \omega_{2},$$
  

$$F_{MD_{1}} = F_{MD_{2}} = f_{0} * D_{w} * B * \sigma_{w}^{2}.$$

The variable drift force spectrum at zero frequency is

$$\begin{split} \mathbf{S}_{F}(0) &\approx 2^{*}(\mathbf{D}_{w}^{*}\mathbf{B})^{2} \int \{f(\omega)^{*}\mathbf{S}(\omega)\}^{2} d\omega, \\ \mathbf{S}_{F_{1}}(0) &\approx 2^{*}(\mathbf{D}_{w}^{*}\mathbf{B}^{*}f_{0}^{*}\mathbf{S}_{1})^{2}^{*} \Delta \omega_{1} = 2^{*}(\mathbf{D}_{w}^{*}\mathbf{B}^{*}f_{0}^{*}\boldsymbol{\sigma}_{w}^{2})^{2} / \Delta \omega_{1}, \\ \mathbf{S}_{F_{2}}(0) &\approx 2^{*}(\mathbf{D}_{w}^{*}\mathbf{B}^{*}f_{0}^{*}\mathbf{S}_{2})^{2}^{*} \Delta \omega_{2} = 2^{*}(\mathbf{D}_{w}^{*}\mathbf{B}^{*}f_{0}^{*}\boldsymbol{\sigma}_{w}^{2})^{2} / \Delta \omega_{2}. \end{split}$$

In particular, the ratio  $S_{F1}(0)/S_{F2}(0)$  can be written simply

$$S_{F_1}(0)/S_{F_2}(0) = \Delta \omega_2 / \Delta \omega_1 = S_1 / S_2$$
.

Note that wave damping of these systems depends essentially on the *mean* drift force, which is the same for these two spectra so that *total* dimensionless system damping  $\Gamma$ , *including* wave-drift contribution, is the same for the two cases (damping in percent of critical = 50\* $\Gamma$ ).

The motion variance of a lightly damped linear spring-mass system with spring constant k, natural frequency  $\omega_0$  and dimensionless damping coefficient  $\Gamma$  is

$$\sigma_{x}^{2} = E[x^{2}] = \pi^{*}\omega_{0}^{*}S_{F}(\omega_{0})/(2\Gamma k^{2})$$

We can as usual approximate  $S_F(\omega_0) \approx S_F(0)$  so the ratio  $\sigma_{x_1} / \sigma_{x_2}$  of rms motions becomes

$$\sigma_{x_1}/\sigma_{x_2} = \sqrt{\{S_{F_1}(0)/S_{F_2}(0)\}} = \sqrt{\{\Delta\omega_2/\Delta\omega_1\}} = \sqrt{\{S_1/S_2\}}$$

Thus, all characteristic motions (rms, most probable peak, etc.) for a linear mooring system will increase roughly in proportion to the square root of the spectral height for constant spectral energy.

*Note:* Since Bretschneider and "Mean" JONSWAP spectra with the same energy and peak period differ in spectral height by about a factor of 2 for typically occurring peak periods, JONSWAP low-frequency vessel motions will be roughly 1.4 times larger than Bretschneider motions for a *linearly moored* vessel, other things being equal.

